

A New Cosmic-Ray-driven Instability

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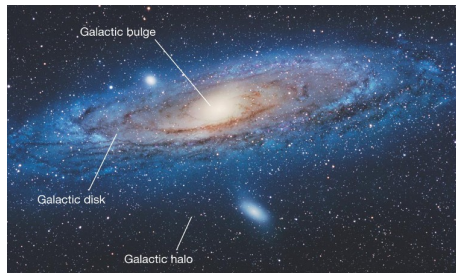
37th International Cosmic Ray Conference

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How do CR couple to background plasma

Estimating time spent by CRs in galactic disk

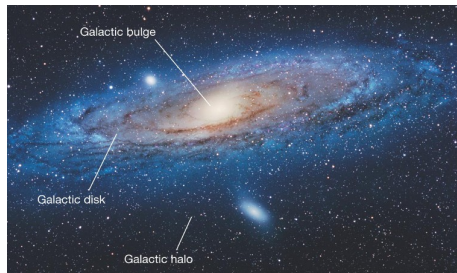
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GeV CR aligned with the field
travel with $c/3$



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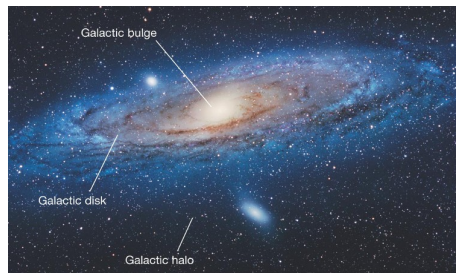
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 $h = 100 - 200 pc$
 $\Rightarrow t_{\text{conf}} \lesssim 10^3$ years.



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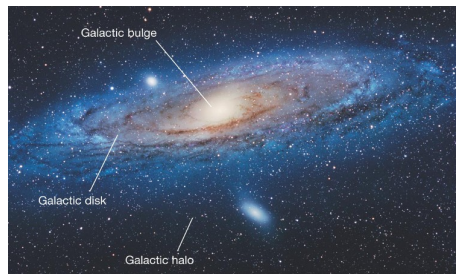
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- ISM:
 $t_{\text{inst}} \sim 10 - 20$ hrs,
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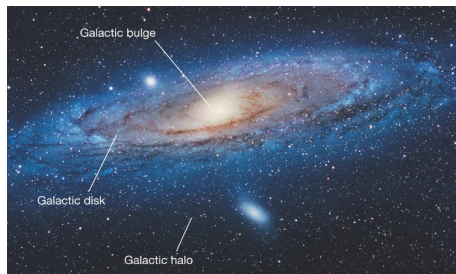
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$$t_{\text{coll}}(i, i) \sim \text{year}$$

\Rightarrow **CR strongly couple by scattering on magnetic field irregularity**

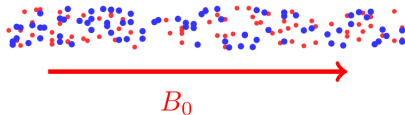


Plan for the talk:

- How GeV cosmic rays couple strongly; plasma instabilities.
- Applications of the new instability
 - ① Electron injection at non-relativistic magnetized shocks.
 - ② CR dynamical impacts on galactic scales (see also Timon Thomas talk).

Electron-ion magnetized plasma

background
plasma

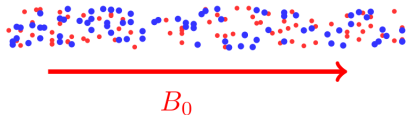


Waves along B_0 :

- Electrostatic

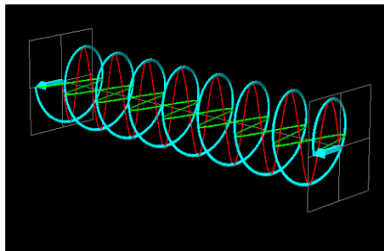
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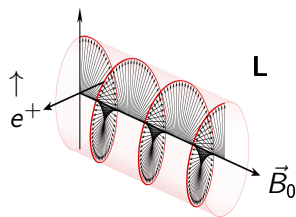
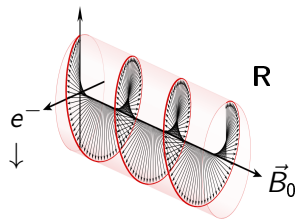
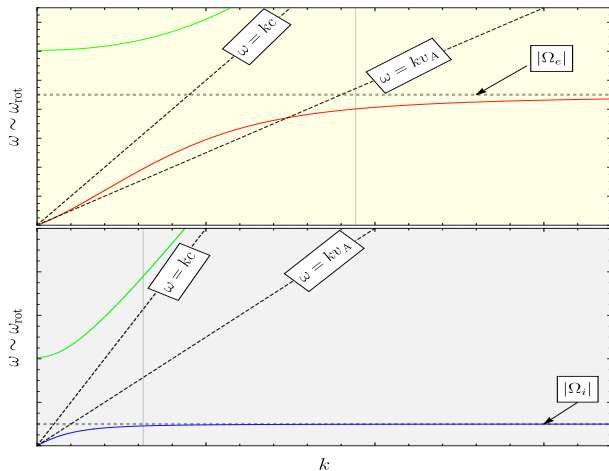


Waves along B_0 :

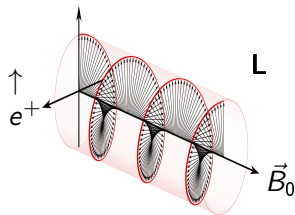
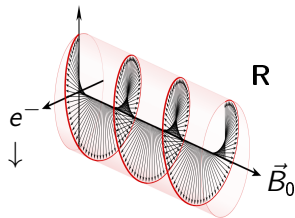
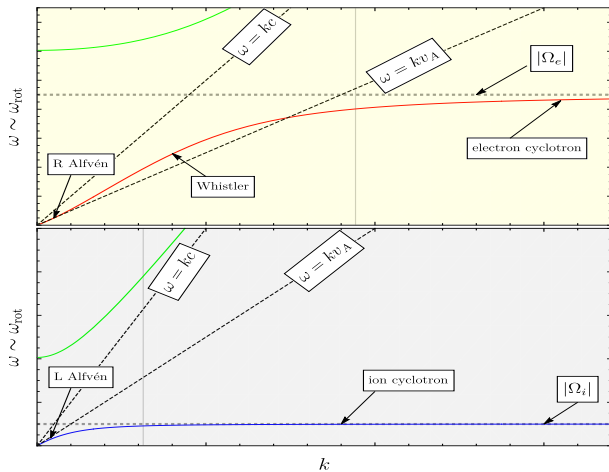
- Electrostatic
- Electromagnetic
Circularly (R & L)
polarized waves



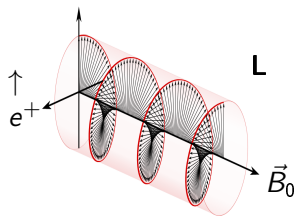
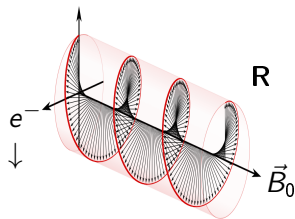
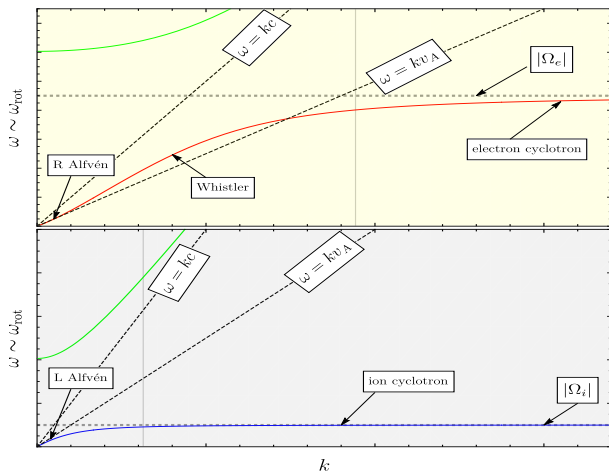
electron-ion magnetized plasma



electron-ion magnetized plasma



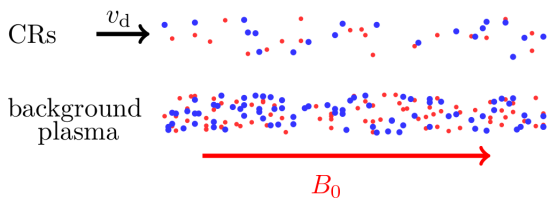
electron-ion magnetized plasma



R Alfvén wave: compressional with $v_{\text{ph}} \gtrsim v_A$

L Alfvén wave: shear with $v_{\text{ph}} \lesssim v_A$

CR driven instability: Gyrotropic CRs



Shalaby+2020; ApJ 908 206

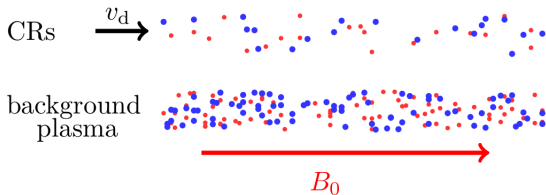
Dispersion relation ($\Omega_{e,0} = -m_i/m_e \Omega_{i,0}$):

$$\frac{k^2 c^2}{\omega^2} - 1 = \frac{\omega_i^2}{\omega(-\omega \pm \Omega_{i,0})} + \frac{\omega_e^2}{\omega(-\omega \pm \Omega_{e,0})} \quad \leftarrow \text{Background}$$

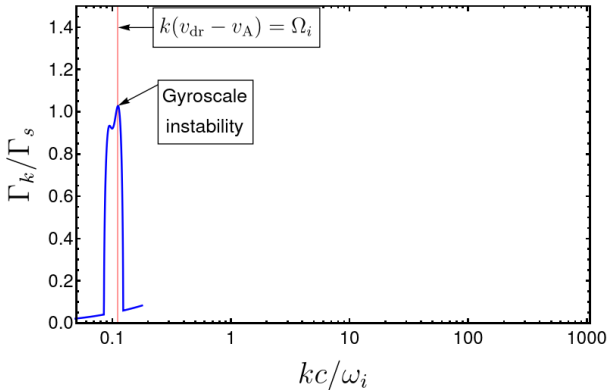
$$\text{CRe} \Rightarrow + \frac{\alpha \omega_e^2}{\gamma_e \omega^2} \left\{ \frac{\omega - k v_{dr}}{k v_{dr} - \omega \mp \Omega_{e,0}/\gamma_e} \right\}$$

$$\text{CRi} \Rightarrow + \frac{\alpha \omega_i^2}{\gamma_i \omega^2} \left\{ \frac{\omega - k v_{dr}}{k v_{dr} - \omega \pm \Omega_i} - \frac{v_{\perp}^2 (k^2 c^2 - \omega^2) / c^2}{2 (k v_{dr} - \omega \pm \Omega_i)^2} \right\}$$

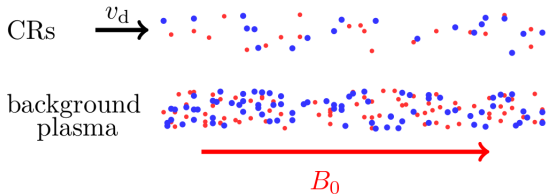
CR driven instability: Gyrotropic CRs



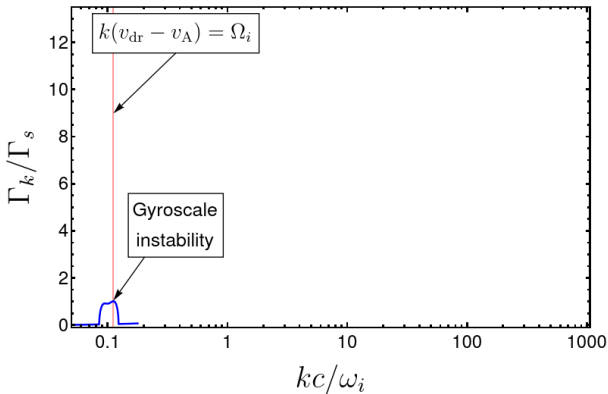
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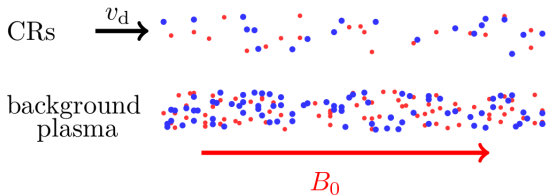
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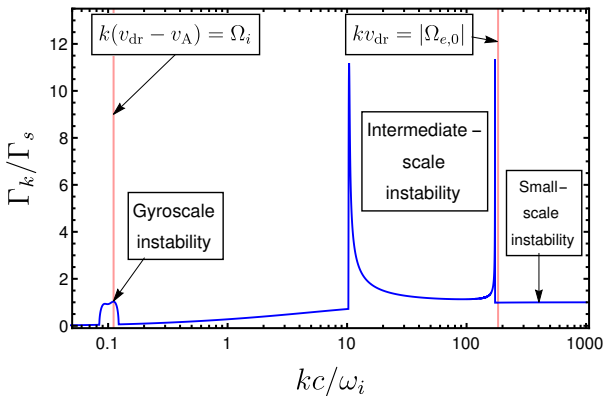
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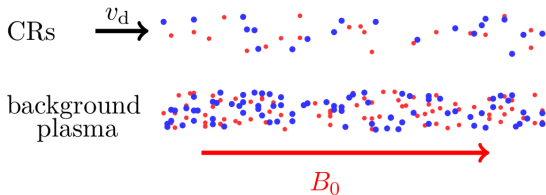
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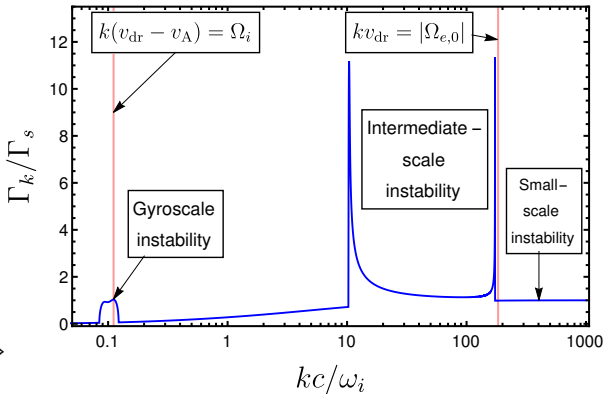
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gyroscale: Alfvén Waves

$$\frac{kc}{\omega_i} = \frac{1}{(v_{dr}/v_A) - 1}$$

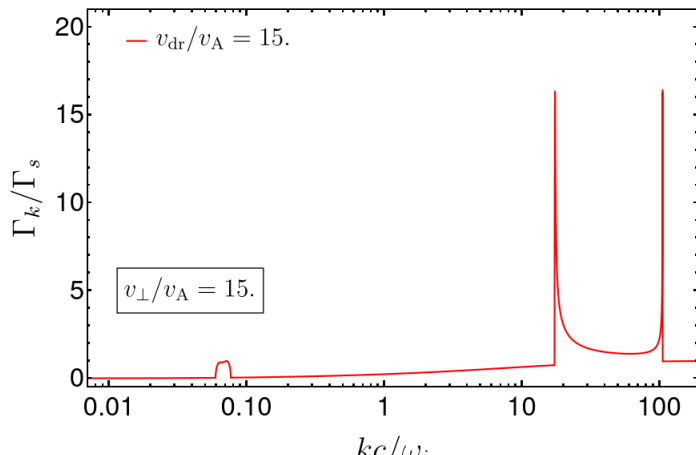
intermediate-scale:
ion-cyclotron waves
(2-peaks)

$$\frac{kc}{\omega_i} \sim \left\{ \frac{v_{dr}}{v_A}, \frac{m_r v_A}{v_{dr}} - \frac{v_{dr}}{v_A} \right\}$$



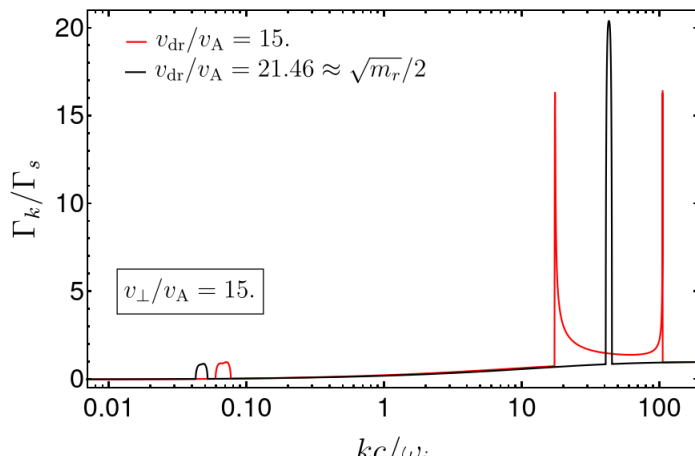
Intermediate-scale: two peaks

$$\frac{kc}{\omega_i} \sim \left\{ \frac{v_{dr}}{v_A}, \frac{m_r v_A}{v_{dr}} - \frac{v_{dr}}{v_A} \right\} \Rightarrow \text{merge} \Rightarrow \frac{v_{dr}}{v_A} = \sqrt{m_r}/2$$



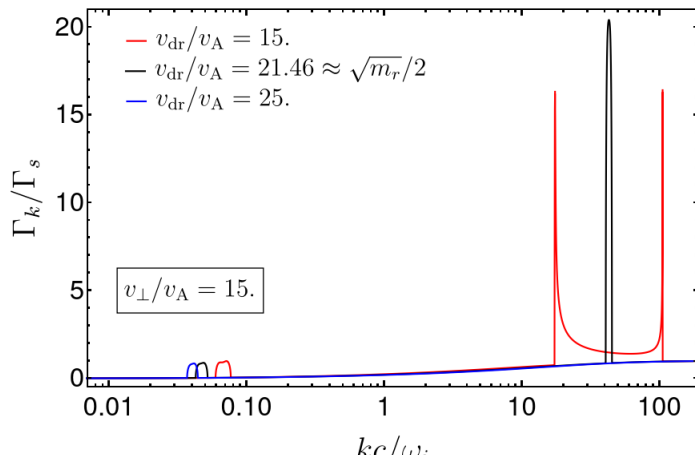
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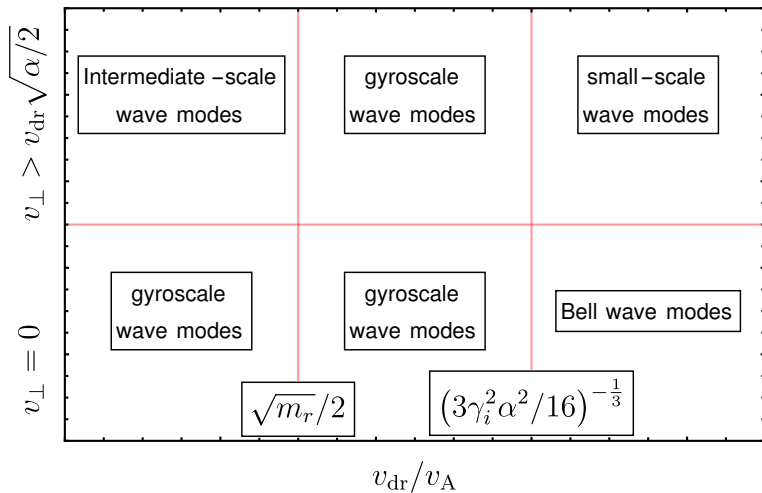


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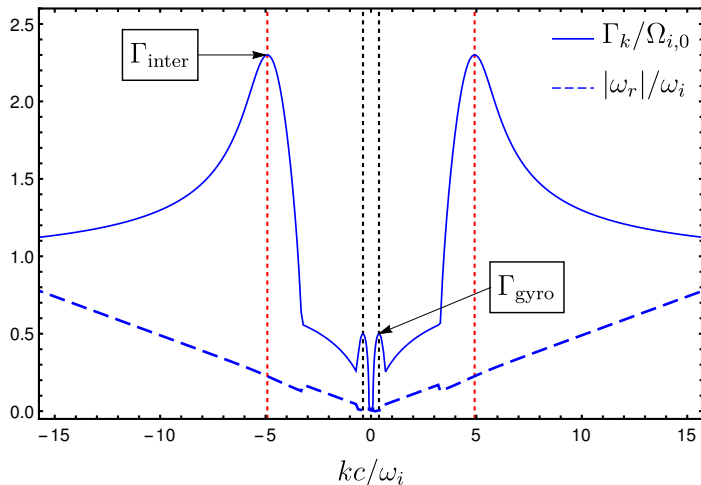
Classification



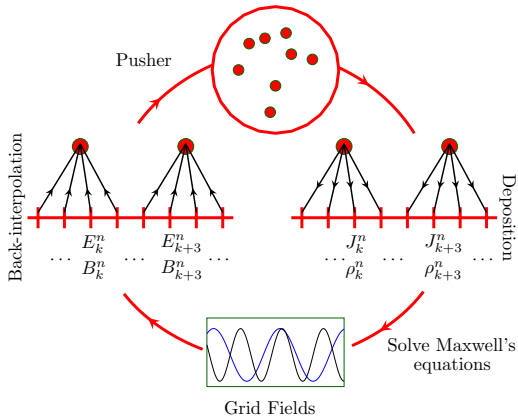
Kinetic simulation for instabilities

Kinetic simulation using Particle-in-Cell:

$v_A = 0.01c$, $m_i/m_e = 1836$, $v_{\text{dr},0} = 5v_A$, $v_{\perp,0} = 13v_A \Rightarrow \theta_0 \sim 70^\circ$



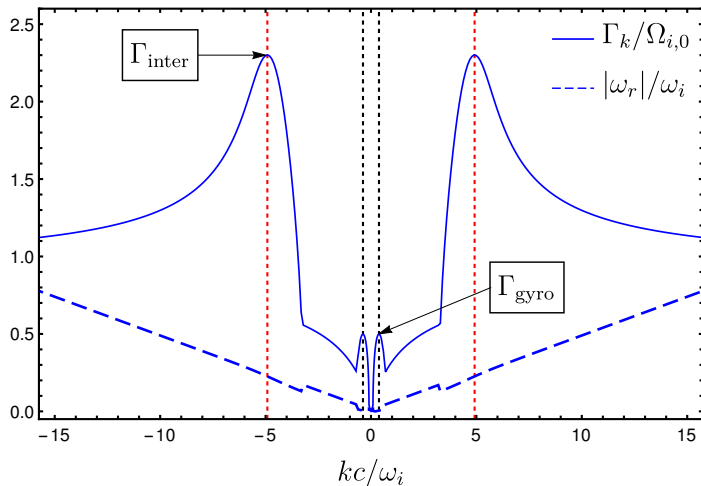
Particle-in-cell algorithm



MS+ 2017, ApJ 841 52

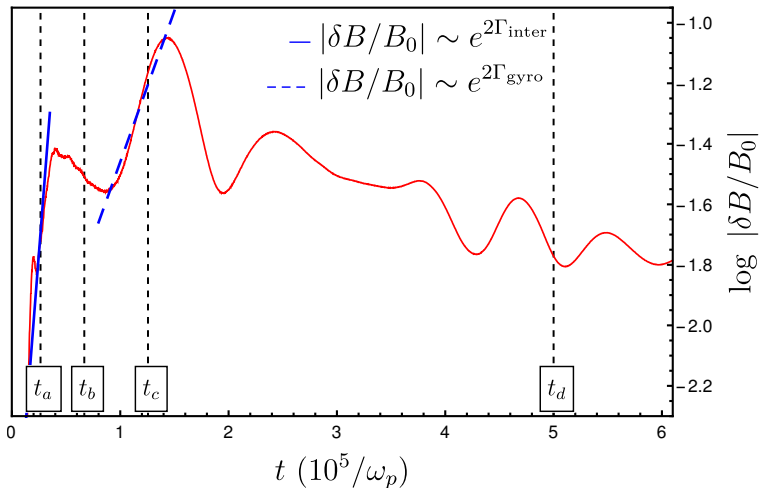
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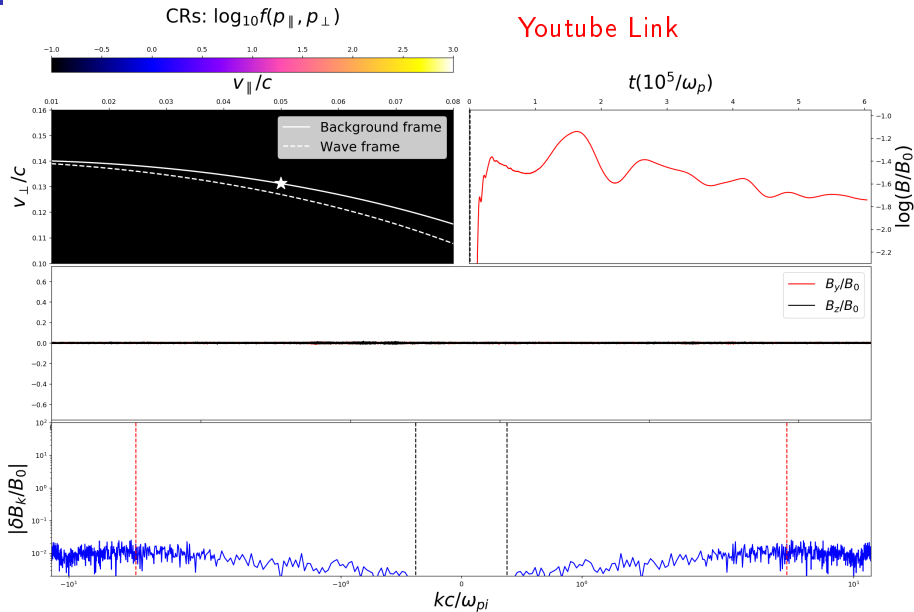


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Kinetic simulation for instabilities



To summarize:

- New instabilities with much faster growth rate if $v_{dr}/v_A < \sqrt{m_r}/2$
- **Only** 2 successful full-kinetic simulations of gyroscale instability
 - ① Holcomb+2019 @ $v_A = 10^{-1}c$, $m_r = 100$ (energy error $\sim 300\epsilon_{cr}$)
 - ② MS+2020 @ $v_A = 10^{-2}c$, $m_r = 1836$ (energy error $\sim 0.002\epsilon_{cr}$)

Both report no-full isotropization in general

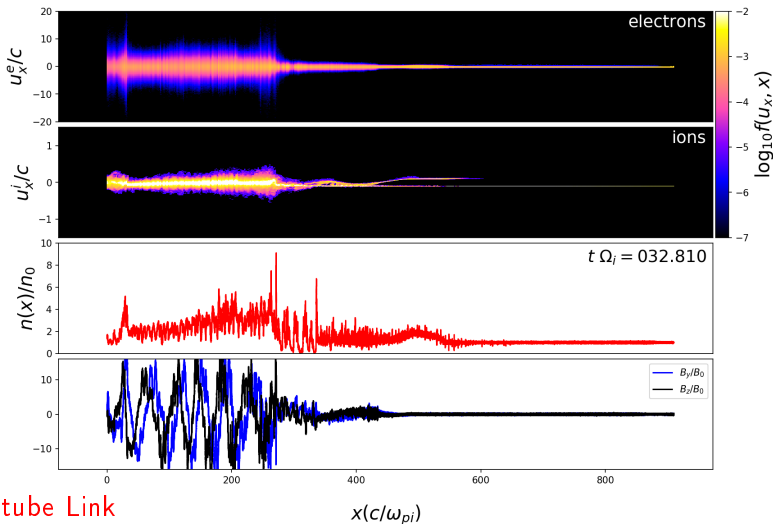
Next: applications

electron injection Problem:

- electrons: $r_e = (m_e/m_i)r_i$.
- electrons can not scatter at shock front
- Intermediate-scale instability provide large-amplitude magnetic perturbation at sub ion-gyroscale
⇒ a solution?

Acceleration at non-relativistic shocks

CD rest frame; $m_i/m_e=1836$; $v_u = -0.1c$; $v_A = 0.00625c$.



[Youtube Link](#)

In the self-confinement picture

$$\frac{d\varepsilon_c}{dt} + \nabla \cdot \left[\vec{W}(\varepsilon_c + P_c) - \kappa \cdot \nabla \varepsilon_c \right] = \vec{W} \cdot \nabla P_c, \quad (1)$$

$$\vec{W} \cdot \nabla P_c = -2 \int d\omega dk \Gamma(\omega, k) I(\omega, k), \quad (2)$$

$$|\kappa| \sim \kappa_{\parallel} \sim \frac{c^2}{2} \left\langle \frac{1 - \mu^2}{\nu_+ + \nu_-} \right\rangle \quad (3)$$

\vec{W} is the effective streaming speed of CRs

New instability: higher linear growth rate

- larger pressure gradient: ∇P_c
- larger scattering rate \Rightarrow lower diffusion coefficient
- Very low Ion-neutral damping rate (10^6 smaller) \Rightarrow mechanism for efficient coupling of MeV CRs to partially ionized plasma, e.g., MC

Conclusions and thoughts

- CR strongly couple via kinetic instabilities
- New instability:
 - much higher rate \Rightarrow new CR transport
 - Can't be suppressed by ion-neutral friction (damping) \Rightarrow potential role in the ionization of molecular clouds by MeV CRs.
- CR impact/regulate galactic outflows and ISM chemistry
- CR transport mode strongly impact CGM gas and magnetic field distribution

Thank you for your attention

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