



# Combined analysis of AMS-02 secondary-to-primary ratios: universality of cosmic-ray propagation and consistency of nuclear cross sections

M. Vecchi, E. F. Bueno, L. Derome, Y. Génolini, and D. Maurin



university of  
 groningen

# Scientific goals

- We study whether F/Si data recently published by AMS-02 [[Aguilar et al Phys.Rev.Lett. 126 \(2021\) 8](#)] can be reproduced by the same propagation models which give a best fit of lighter secondary-to-primary ratios, as derived in [Weinrich et al, A&A 639, A131 \(2020\)](#)
- We follow the methodology described in [Derome et al, A&A 627 \(2019\) A158](#)
- We investigate whether data allow for primary F component

**NB:** CR fluorine is purely composed of (stable)  $^{19}\text{F}$

# Cosmic-ray nuclei

**Primaries** are produced and accelerated at the sources.

**Secondaries** are produced by the collisions of **primaries** with the **interstellar medium (ISM)**.

**Primaries (H, O, Si, ...)**

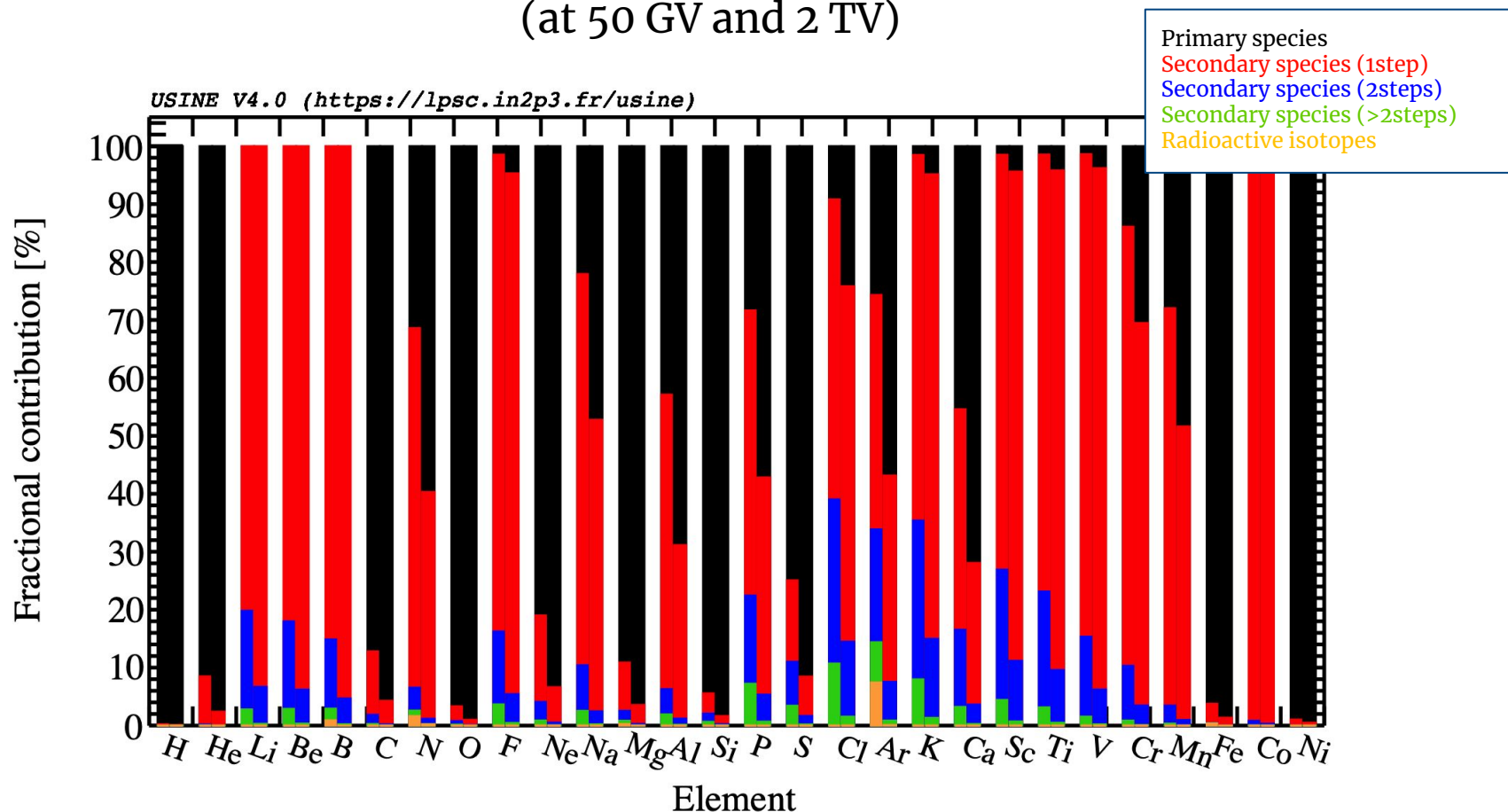
A diagram illustrating the propagation of cosmic ray nuclei in the Galaxy. A wavy orange line represents the path of primary cosmic rays, starting from the bottom left and moving towards the center. A red starburst marks a collision point where a primary nucleus interacts with the interstellar medium. From this point, a wavy green line represents the path of secondary cosmic rays, which are produced by the collision and travel towards the top right. The background is a starry field with a prominent spiral galaxy structure.

**Secondaries (D, B, E, ...)**

Secondary-to-primary flux ratios, such as B/C or F/Si, are key observables to constrain the propagation processes in the Galaxy.

# Secondary CR production

Relative contributions per production process for elemental fluxes  
(at 50 GV and 2 TV)



The species with the highest primary content include H, O, Si, and Fe (black), while Li, Be, B, F, and Cl to V have the highest secondary component from both single (red) and multi-step production (blue and green).

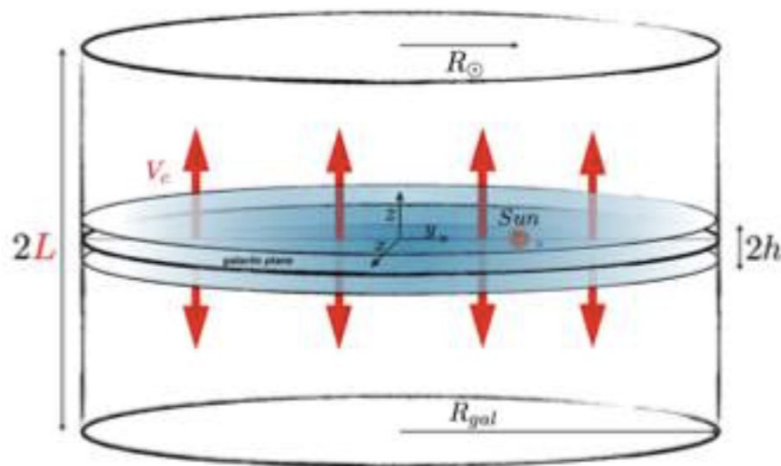
# Cosmic-ray transport in the Galaxy

$$\begin{aligned}
 -\vec{\nabla}_{\mathbf{x}} \left\{ K(E) \vec{\nabla}_{\mathbf{x}} \psi_{\alpha} - \vec{V}_c \psi_{\alpha} \right\} + \frac{\partial}{\partial E} \left\{ b_{\text{tot}}(E) \psi_{\alpha} - \beta^2 K_{pp} \frac{\partial \psi_{\alpha}}{\partial E} \right\} \\
 + \sigma_{\alpha} v_{\alpha} n_{\text{ism}} \psi_{\alpha} + \Gamma_{\alpha} \psi_{\alpha} = \underbrace{q_{\alpha}}_{\text{source}} + \sum_{\beta} \left\{ \sigma_{\beta \rightarrow \alpha} v_{\beta} n_{\text{ism}} + \Gamma_{\beta \rightarrow \alpha} \right\} \psi_{\beta}
 \end{aligned}$$

$K(E)$ : A two-break diffusion coefficient is used

Gérolini et al PRL 119, 241101 (2017), Gérolini et al Phys.Rev. D99 (2019)

$q_{\alpha}$ : A universal single power-law is used for the source term.



1D model and semi-analytic approach with the USINE code

[Maurin CPC 247 (2020) 106942, <https://dmaurin.gitlab.io/USINE/>]

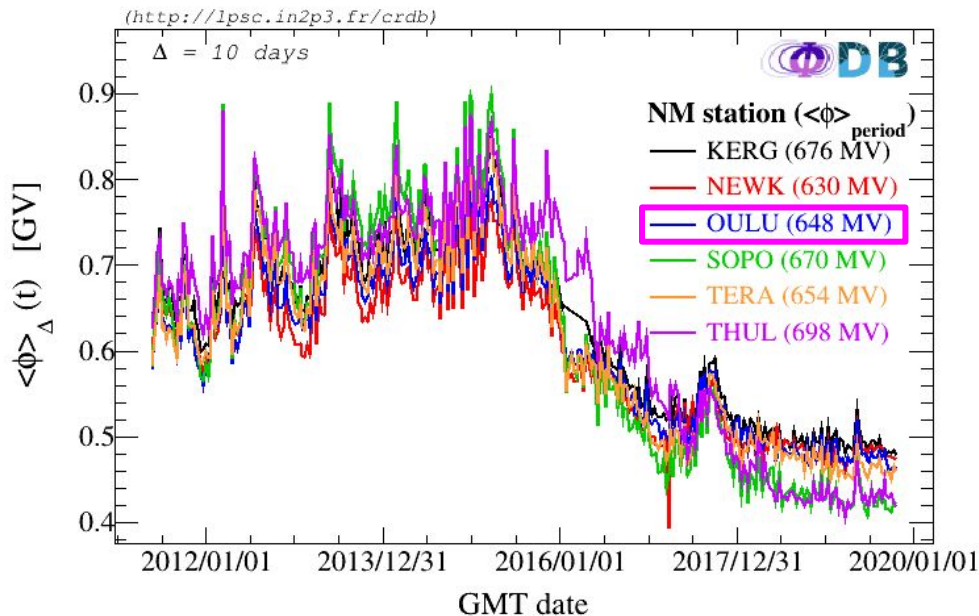
# Cosmic-ray transport in the Galaxy

$$-\vec{\nabla}_{\mathbf{x}} \left\{ K(E) \vec{\nabla}_{\mathbf{x}} \psi_{\alpha} - \vec{V}_c \psi_{\alpha} \right\} + \frac{\partial}{\partial E} \left\{ b_{\text{tot}}(E) \psi_{\alpha} - \beta^2 K_{pp} \frac{\partial \psi_{\alpha}}{\partial E} \right\} \\ + \sigma_{\alpha} v_{\alpha} n_{\text{ism}} \psi_{\alpha} + \Gamma_{\alpha} \psi_{\alpha} = q_{\alpha} + \sum_{\beta} \left\{ \sigma_{\beta \rightarrow \alpha} v_{\beta} n_{\text{ism}} + \Gamma_{\beta \rightarrow \alpha} \right\} \psi_{\beta}$$

- This equation couples about a hundred CR species (for  $Z < 30$ ) over a nuclear network of more than a thousand reactions.
- To solve this diagonal matrix of equations, we start with the heaviest nucleus, which is always assumed to be a primary species, and then proceed down to the lightest one.
- We use the SLIM propagation scenario [[Génolini et al Phys.Rev. D99 \(2019\)](#)], with 4 free propagation parameters ( $K_0, \delta, R_l, \delta_l$ )

# Methodology

- In order to reduce biases in the transport parameter determination, it is crucial to use nuisance parameters for the nuclear production cross sections, and a covariance matrix for the data systematic uncertainties, as described in [Derome et al, A&A 627 \(2019\) A158](#)
- The force-field approximation is used to compute the top-of-atmosphere (TOA) fluxes, using the Fisk potential as a nuisance parameter.

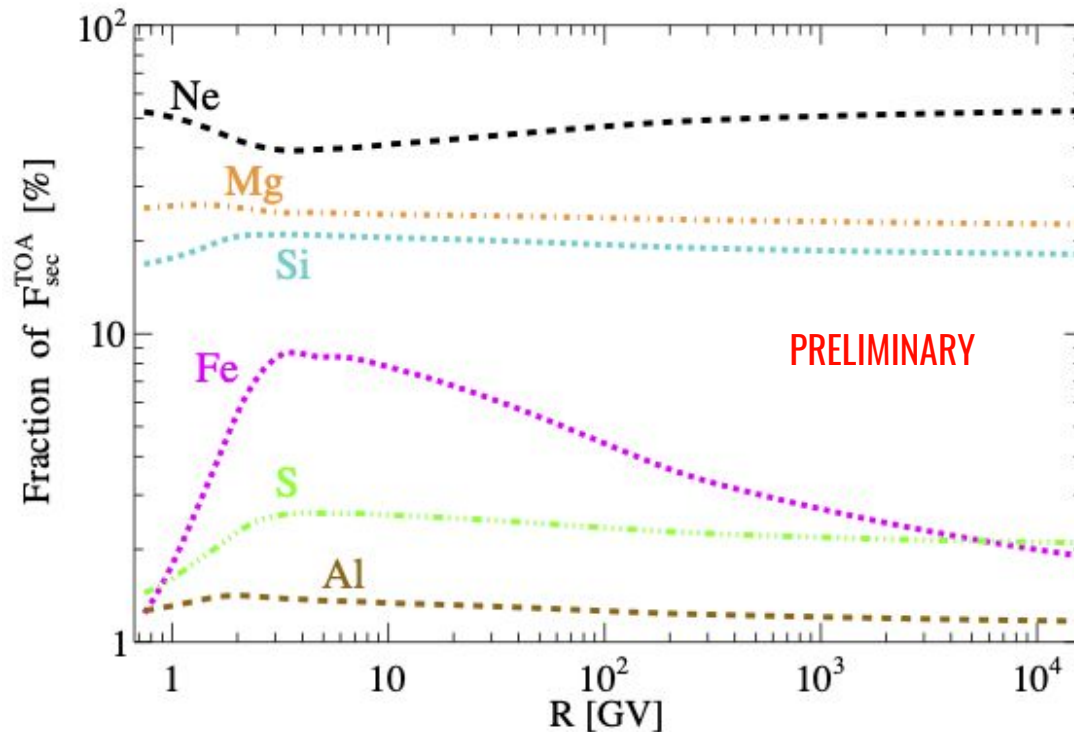


$\langle \Phi \rangle$  from <https://lpsc.in2p3.fr/crdb/>  
based on [Ghelfi et al., AdSR 60, 833 \(2017\)](#)

- The TOA fluxes are compared to the data using a  $\chi^2$  minimization procedure

# Progenitors of CR fluorine

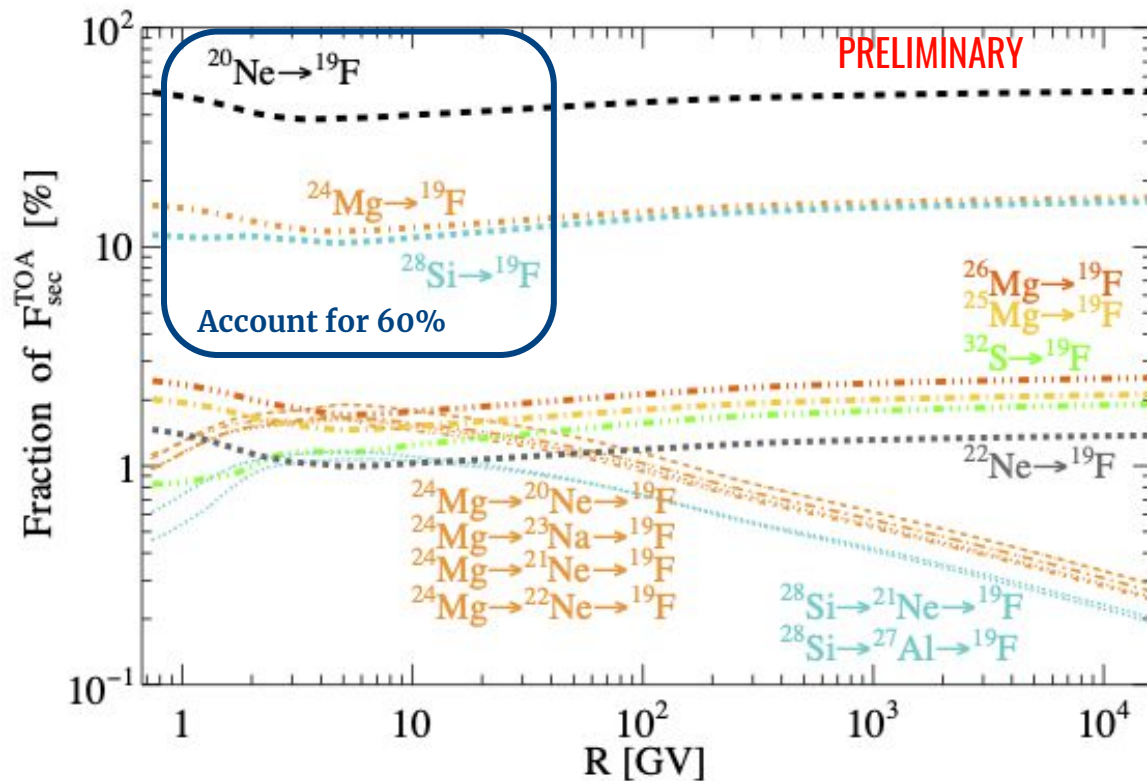
Following the methodology described in [Génolini et al \*Phys.Rev.C\* 98 \(2018\) 3, 034611](#)



Ne, Mg, Si and Fe are the main progenitors of F.



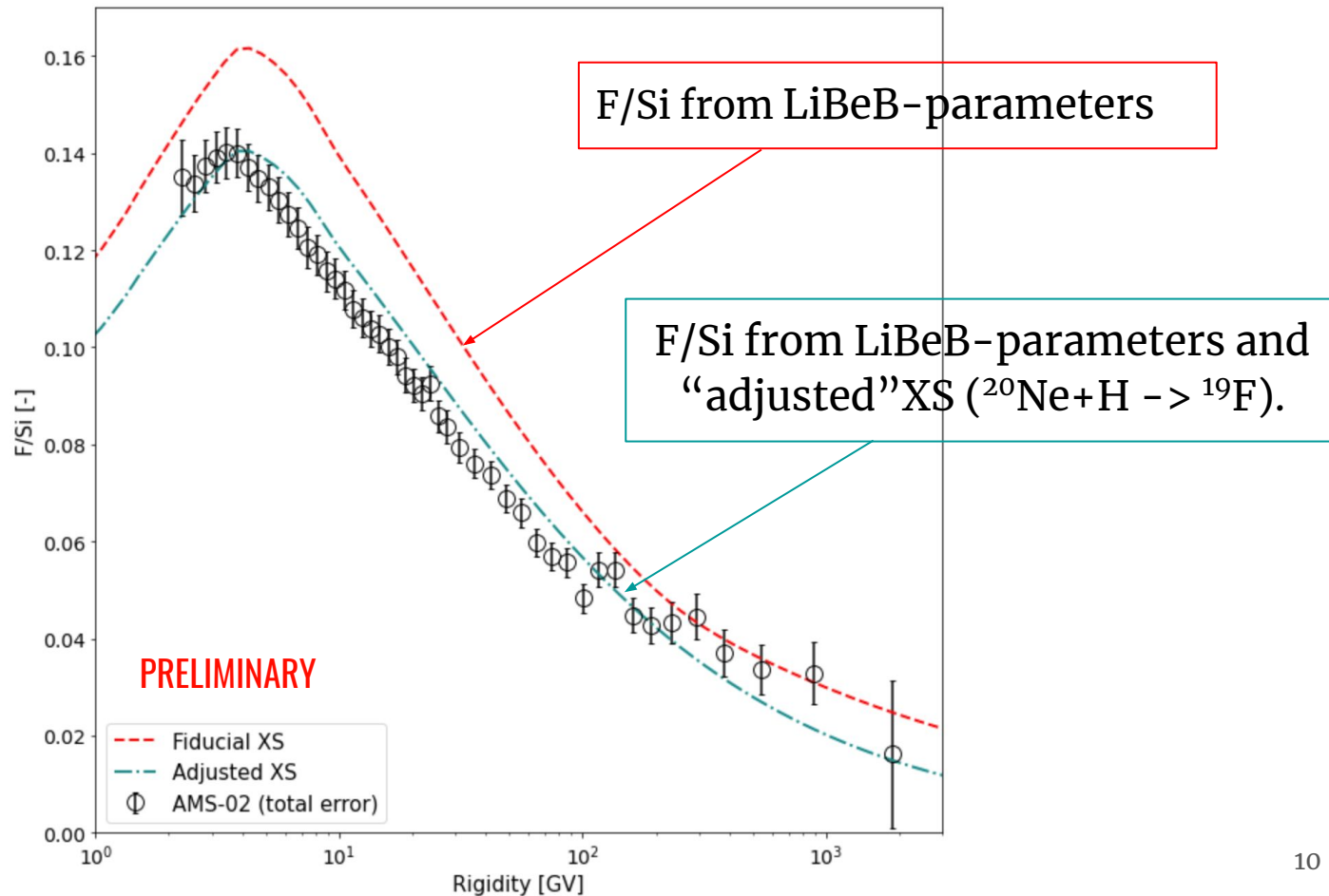
# Dominant processes producing CR fluorine



- We have identified **35 channels** whose individual contribution is in the range [0.1-2%] and contribute to  $\sim 22\%$  of the total.
- While the ranking of the dominant channels is a robust prediction, the individual numbers are subject to uncertainties due to the cross section and propagation parameters.

# Results

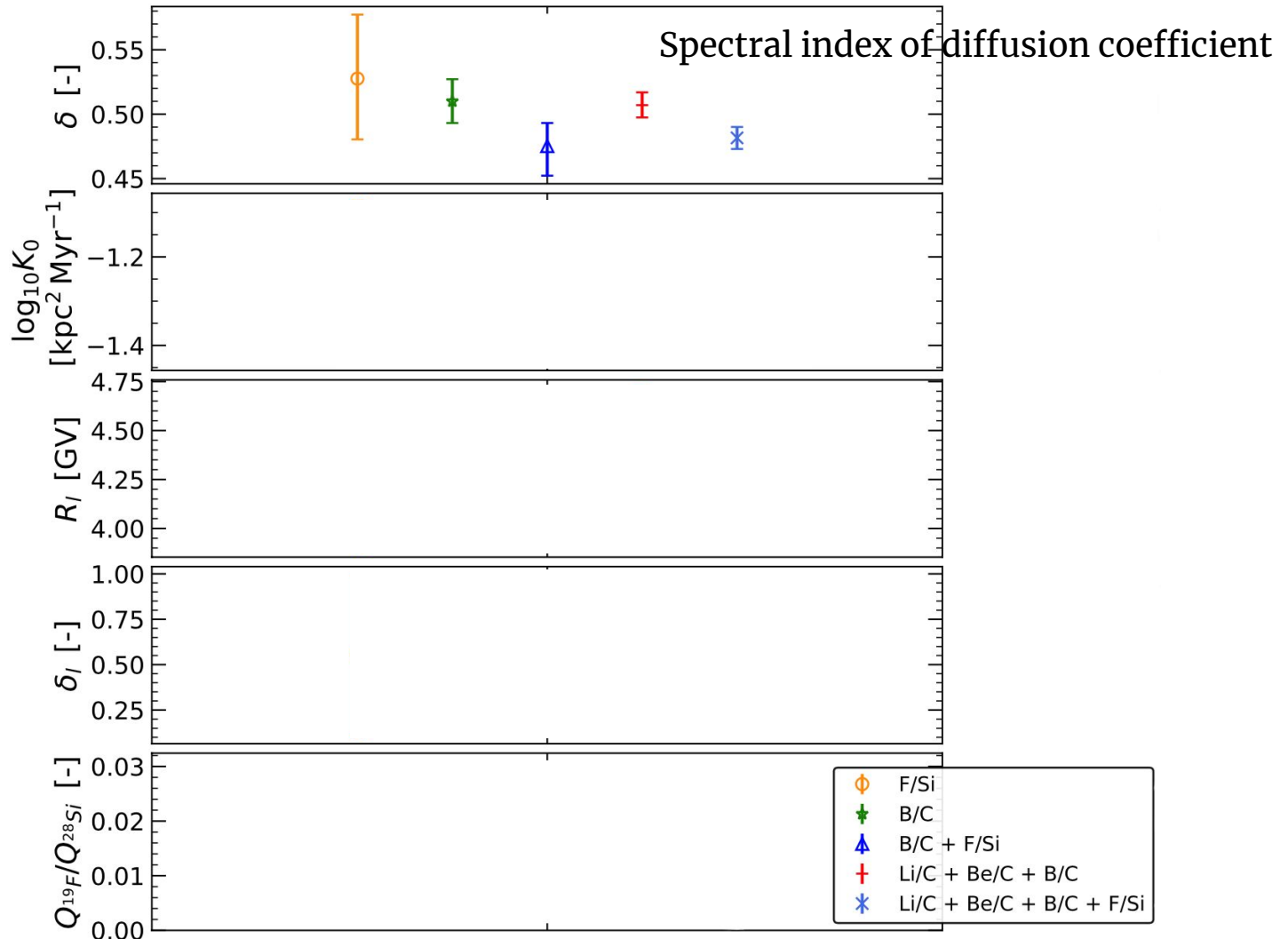
- The model tuned on light nuclei AMS-02 data [Weinrich et al, A&A 639, A131 (2020)] overshoots F/Si data by 10% - 15%, similar to M. Boschini et al, arXiv:2106.01626
- o(15%) reduction in the XS normalization to recover the difference, in the range of expected XS uncertainties



# Results: propagation parameters (and more)

Effective  
diffusion  
coefficient

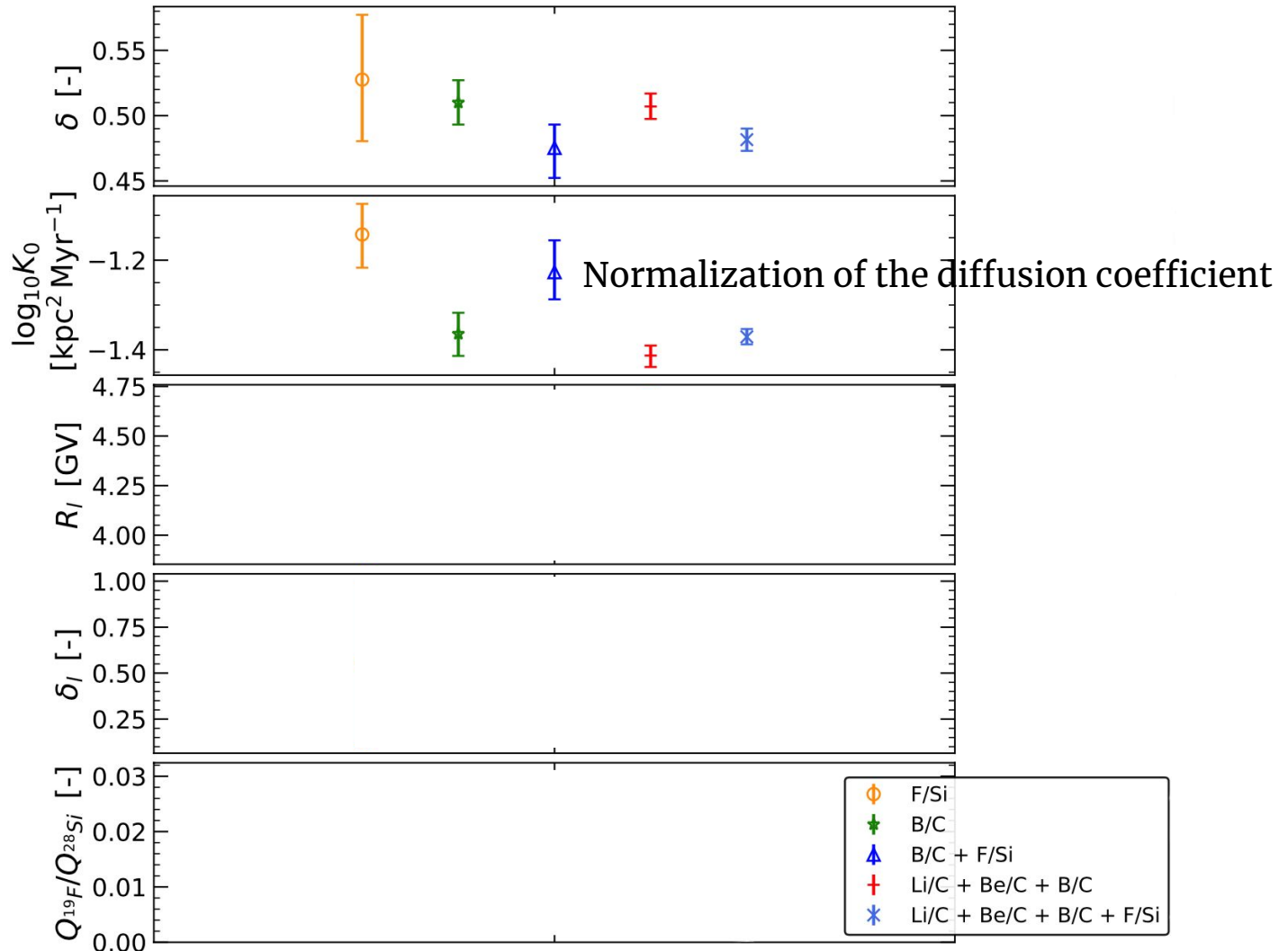
$$K(R) = \beta^\eta K_0 \left\{ 1 + \left( \frac{R}{R_1} \right)^{\frac{\delta_1 - \delta}{s_1}} \right\}^{s_1} \left\{ \frac{R}{R_0 = 1 \text{ GV}} \right\}^{\delta} \left\{ 1 + \left( \frac{R}{R_h} \right)^{\frac{\delta - \delta_h}{s_h}} \right\}^{-s_h}$$



# Results: propagation parameters (and more)

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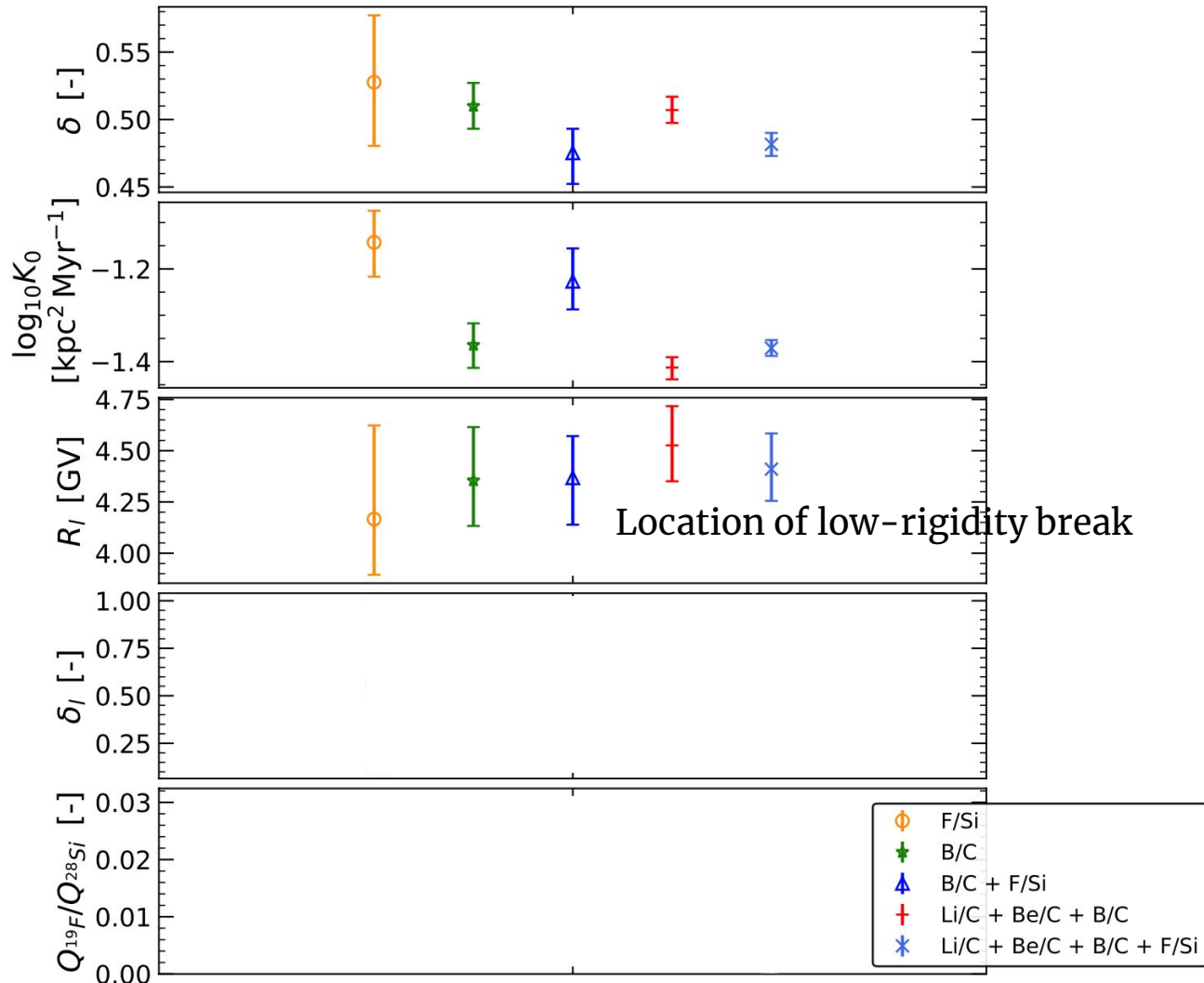
$$K(R) = \beta^n K_0 \left\{ 1 + \left( \frac{R}{R_1} \right)^{\frac{\delta_1 - \delta}{s_1}} \right\}^{s_1} \left\{ \frac{R}{R_0 = 1 \text{ GV}} \right\}^\delta \left\{ 1 + \left( \frac{R}{R_h} \right)^{\frac{\delta - \delta_h}{s_h}} \right\}^{-s_h}$$



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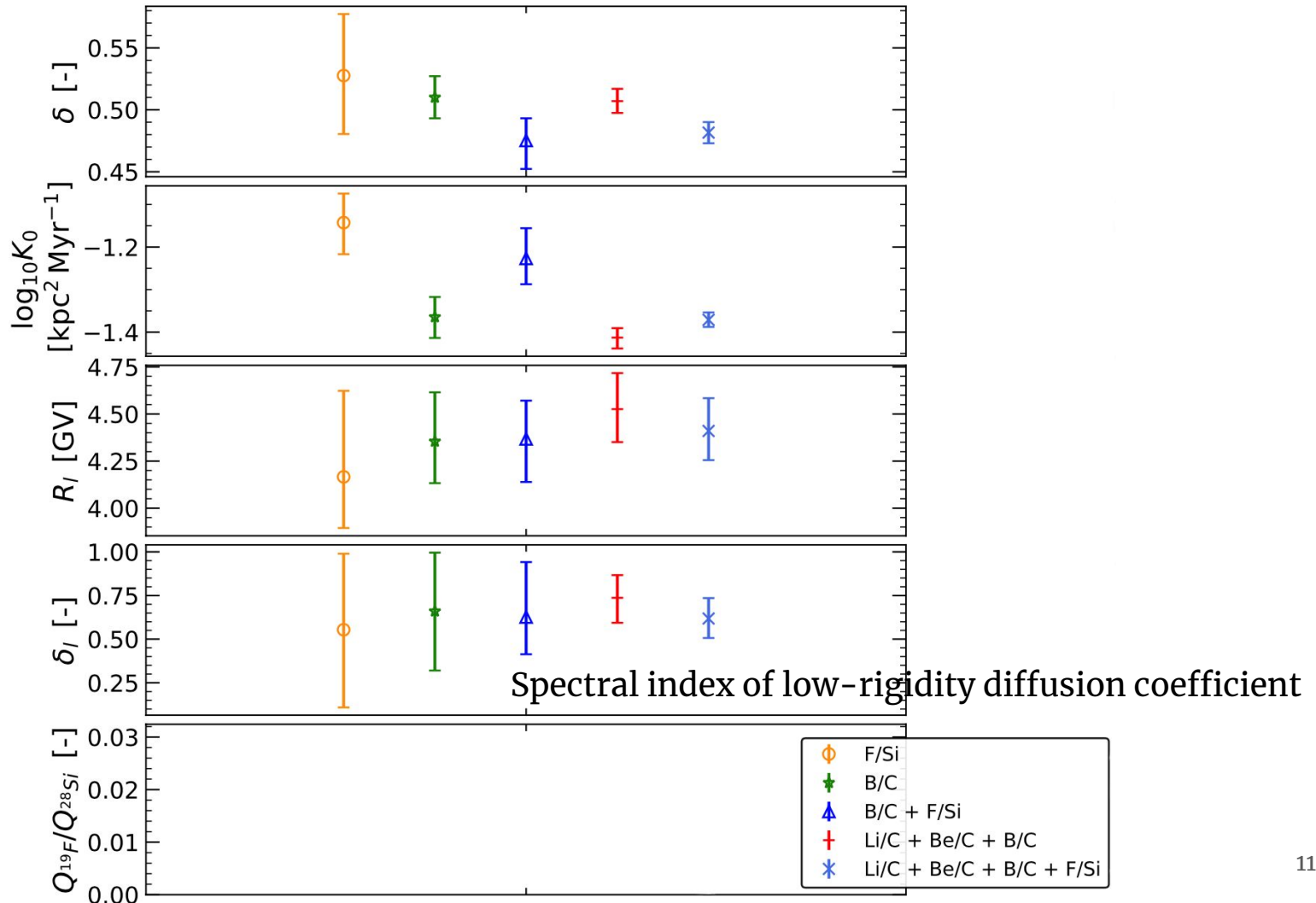
$$K(R) = \beta^\eta K_0 \left\{ 1 + \left( \frac{R}{R_l} \right)^{\frac{\delta_l - \delta}{s_l}} \right\}^{s_l} \left\{ \frac{R}{R_0 = 1 \text{ GV}} \right\}^\delta \left\{ 1 + \left( \frac{R}{R_h} \right)^{\frac{\delta - \delta_h}{s_h}} \right\}^{-s_h}$$



# Results: propagation parameters (and more)

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diffusion  
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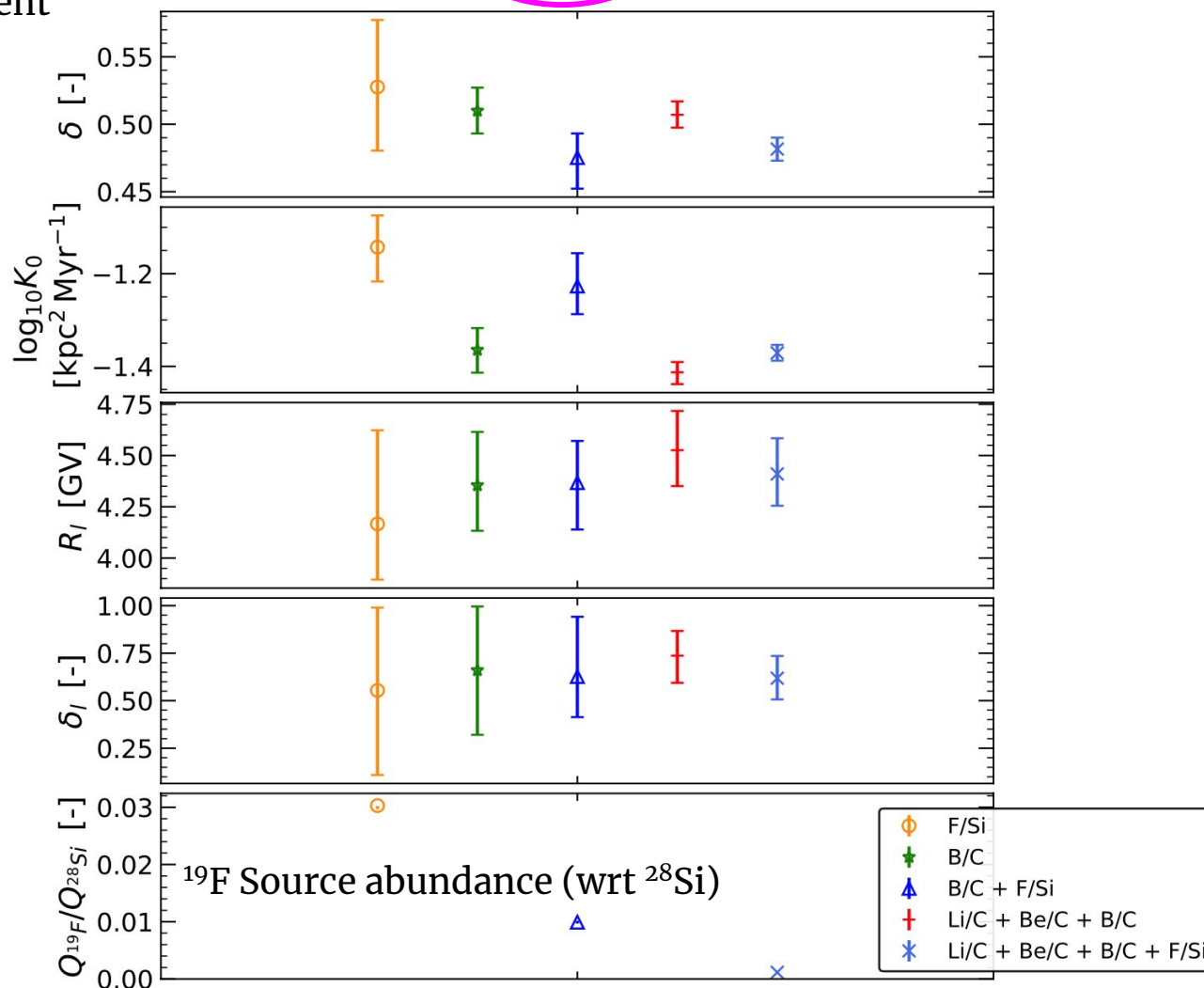
$$K(R) = \beta^\eta K_0 \left\{ 1 + \left( \frac{R}{R_l} \right)^{\frac{\delta_l - \delta}{s_l}} \right\}^{s_l} \left\{ \frac{R}{R_0 = 1 \text{ GV}} \right\}^\delta \left\{ 1 + \left( \frac{R}{R_h} \right)^{\frac{\delta - \delta_h}{s_h}} \right\}^{-s_h}$$



# Results: a F primary component ?

Despite  $^{19}\text{F}$  being mostly secondary, we study the effect of a primary component

$$Q(R) = Q_{19\text{F}} \beta^{\eta_S} R^{-\alpha}$$



# Conclusions

- Using the propagation parameters which give a best fit of lighter secondary-to-primary ratios, our model overestimates the data by 10% - 15%. However, this difference can be explained by the F production cross-sections uncertainties
- We conclude that all secondary species from Li to F can be explained by the same transport parameters
- Combined analysis of Li/C, Be/C, B/C and F/Si gives an upper limit on the F source abundance
- Next steps:
  - Final check on the x-sections
  - Interpret the limit on the F source abundance